

A unified phenomenological description for the origin of mass for leptons and for the complete baryon octet, including the reported inverse dependence upon the alpha constant

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PACS: 03.65.Sq; 12.39.-x

Keywords: zitterbewegung, mass, leptons, baryons

Abstract

Several authors have reported the dependence of the rest masses of particles upon the inverse of the alpha constant. Barut was able to associate such behavior with magnetic selfenergy effects in the case of leptons. The present author has taken account of such magnetic energy effects phenomenologically, in a way similar to Post's, many years ago. This paper presents the extension of the approach to the full baryon octet, and the inverse dependence with alpha is obtained. The masses of all these particles are shown to be described in terms of magnetodynamic energies considering as a fundamental feature the quantization of magnetic flux inside a zitterbewegung motion "orbit" performed by each particle in consequence of its interaction with the vacuum background.

Introduction

Several authors have reported the dependence of the rest masses of particles upon the inverse of the alpha constant. Barut was able to associate such behavior with magnetic selfenergy effects in the case of leptons[1]. The present author has taken account of such magnetic energy effects phenomenologically[2], in a way similar to Post's, many years ago[3]. This paper presents the extension of the approach to the full baryon octet, and the inverse dependence with alpha is obtained. The masses of all these particles are shown to be described in terms of magnetodynamic energies considering as a fundamental feature the quantization of magnetic flux inside a zitterbewegung motion "orbit" performed by each particle in consequence of its interaction with the vacuum background(Jehle[4] proposed flux quantization inside zitterbewegung orbits of particles as early as 1967).

Our previous work begins with the concept of gauge invariance and consequent flux quantization associated with the zitterbewegung intrinsic motion of fundamental particles. We then associated the magnetodynamic energy of the motion with the rest energy of a particle[2,3]. The main result of such phenomenological analysis was eq. (3) of [2]:

$$\frac{mR^2}{\mu} = \frac{nh}{2\pi ec} \quad (1)$$

In this equation m is mass, R is the range of the vibrational intrinsic motion of the particle, μ is the magnetic moment, n is the number of magnetic flux quanta(admitted as given by the nonrelativistic expression hc/e). The model adopts experimental values for m and μ . For the nucleons R was given by theoretical values calculated by Miller[5], and for the electron (and the muon) this parameter was assumed as equal to the Compton wavelength $\lambda = \hbar/mc$ [6]. Good agreement between model and experiment was obtained for that reduced group of particles.

However, the application of the model to other particles depends on the knowledge of the parameter R . In order to put the model to further test, in the present work we decided to simply try and eliminate the explicit dependence of the model upon R . For the leptons the following expression is known to be valid:

$$\mu = e\lambda/2 \quad (2)$$

Here $\mu = \mu_B$ is the magnetic moment in the case of the electron (μ_B is the Bohr magneton). Therefore, for the leptons and for the members of the baryon octet considered in this work we will *assume* that in (2) $\lambda/\sqrt{2}$ can be directly replaced by R , so that R is eliminated from (1) in favor of μ . It is clear that such possibility associates mass to only two parameters, namely, the number of flux quanta imposed by gauge invariance conditions and the charges of the constituents inside the baryons, and to the inverse of the experimental magnetic moment. As shown below we verify that such proposal is consistent with experiment.

Inserting the definition for R into (1) and using the definition of the fine structure constant alpha, $\alpha = e^2/\hbar c$, we can rewrite (1) in the form:

$$\frac{2c^2\alpha}{ne^3} m = \frac{1}{\mu} \quad (3)$$

It can immediately be noticed that n and μ should in principle be directly associated with each other, which would produce an inverse dependence of m with the alpha constant, as reported in the literature

Application to Leptons and Baryons

A.O.Barut [7,8] proposed an alternative theory for the inner constitution of baryons and mesons, in which the basic pieces would be the individual, *stable* particles, namely the proton p , the electron e^- , and the neutrino ν (ν' will indicate antineutrinos, below), rather than quarks with fractionary charges which do not manifest themselves physically as individual entities. The muon μ^- would also be included and considered responsible for effects usually denoted as “strangeness”. Barut proposed also that the short range strong interactions between such internal constituents would be magnetic in nature. In order to account for the same conservation rules as a model based upon the fractionary quarks does, the constitution of baryons should be as follows[7,8]: proton = $p = (p e^- e^+)$, neutron = $n = (p e^- \nu)$, $\Sigma^- = (p e^- \mu^- \nu)$, $\Sigma^0 = (p \mu^- \nu)$, $\Sigma^+ = (p e^+ \mu^-)$, $\Xi^- = (p \mu^- \mu^- \nu)$, $\Xi^0 = (p \mu^- \mu^- e^+ \nu)$, $\Lambda = (p \mu^- \nu \nu)$. We see that the proton is present in all baryons but is itself a composite particle, supposedly containing an electron and a positron. The analysis below shows that assuming Barut’s ideas are correct, our recently proposed model (which assumes particles inside the baryons can

be considered individually) can also be applied to the full baryon octet with almost perfect accuracy. However, the precise values of n , the number of flux quanta in (3), are actually unknown. The determination of these numbers would require the knowledge of the proper topological properties of each baryon and how to sum individual contributions from its constituents. Relativistic effects if relevant would certainly also have an effect on these numbers, which might even be half-integers. A previous attempt, in a model that also relates particles to zitterbewegung was proposed by Jehle[4], associating particles to the topology of torus knots. Instead of a single n Jehle associates flux quantization to a complicated combination of winding and whirling numbers. However, in the spirit of a phenomenological model, in Table 1 we notice that the magnetic moments for the baryons in the last column are almost perfectly ordered in small numbers of nuclear magnetons. Considering that the magnetic moments should be proportional to the number of flux quanta trapped in the zitterbewegung motion we adopted the following procedure. At least one of these numbers, $n=3$ for the proton, can be justified in simple terms by assuming spin-1/2 quarks as forming the proton. In a previous (unpublished) calculation based upon an average over the three different quarks spins configurations weighted by their Clebsch-Gordan coefficients, we obtain exactly $n=3$ flux quanta for a proton. This is quite close to the magnetic moment of the proton in number of nuclear magnetons. For the other baryons we extrapolate this result by taking for n the integer or half-integer number which is closest to the observed magnetic moment in nuclear magneton units. The results are in Table 1 and we immediately notice that the ratio n/μ is approximately the same for all baryons. The magnetic moment data in the Table are from [9].

Analysis

Figure 1 shows the plot of eq.(3) and the straight solid line would indicate perfect agreement with theory. We observe that Equation (3) describes perfectly well the data available for leptons (triangles) and baryons with the values of n in Table 1.

The influence of topology(introduced through the concepts of flux quantization and gauge invariance) is evident in view of the importance of the empirical sequence of values for n in Table1, and their clear association

with the actual magnetic moment data, which can only be interpreted in such geometrical terms.

There exists a wealth of references in the literature in which scaling laws are proposed based [10-12] on experimental results, to associate mass for all particles with the inverse of α . We see from eq. (3) and Figure 1 that such relation with α indeed is part of our results, since the ratio n/μ is essentially the same for all baryons. In particular, the analysis in ref. [12] might probably be reproduced if the ratio n/μ in (3) is made part of the *free* parameter N in ref. [12].

Conclusions

In resume, this paper has shown that if one properly inserts quantum conditions in a closed-orbit intrinsic motion for the fundamental particles (even in a nonrelativistic limit), in order that gauge invariance is introduced in the treatment, the masses for these particles are directly dependent only upon the inverse of their magnetic moments and upon the number of magnetic flux quanta inside the orbits. This demonstrates the influence of geometrical or topological effects on the problem of mass determination. On the other hand a model like this is only justifiable if baryons are formed by the combination of bonafide stable particles of unitary or null electronic charge, as proposed by Barut, since the introduction of quarks instead would bring in the usual uncertainties related to their formal description in terms of conventional wavefunctions, and so on.

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Table 1: Data utilized in Figure 1. The values of n are chosen as the integer or half-integer numbers that follow as close as possible the (apparent !) sequence in the last column for the baryons, in order to fit theory to data. The magnetic moments are from ref. [9]. One needs to convert mass to grams, magnetic moments to erg/gauss (all CGS units).

part	Rest energy(MeV)	n	(Abs)Magnetic moment(n.m.)
e	0.511	1	1836
muon	105.66	1	8.89
p	938.27	3	2.79
n	939.56	2	1.91
Σ^+	1189	2.5	2.46
Σ^0	1192	1	~ 0.7 (theor.)
Σ^-	1197	1.5	1.16
Ξ^0	1314	1.5	1.25
Ξ^-	1321	1	0.65
Λ	1116	0.5	0.61

Figure 1: Plot of eq. (3). The dotted lines indicate a factor of 2 around the solid line. Triangles are leptons and circles are baryons.

